

Incorporation of spatially extended bodies in MCell (TR&D1 and TR&D2)

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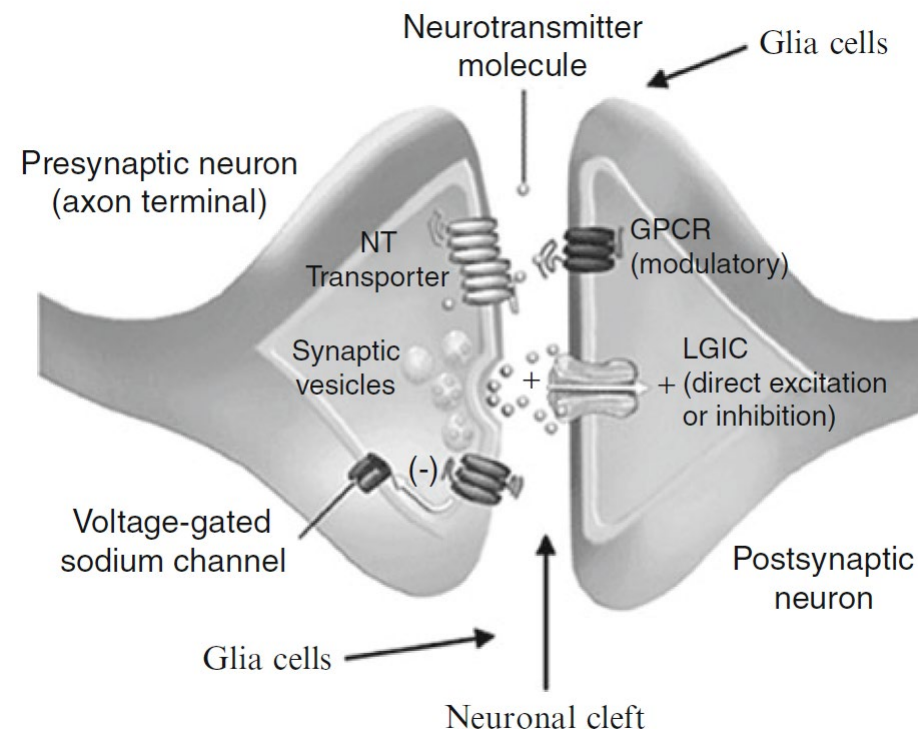
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Current state of MCell:

In MCell[†], reaction-diffusion problems of a set of molecules are modelled as stochastic processes based on Monte-Carlo methods within a region, which can be either a surface of a (pre/post)-synaptic neuron, or the space bounded by those surfaces, the so-called neuronal cleft.

- Each molecule is modelled as a point-particle.
- Diffusion of each particle is governed by an overdamped stochastic dynamics.
- The orientation of the possible binding sites does not play a role (structural information).
- Non-binding particles passes through each other.
- Codes governing the motion of particles in \mathbb{R}^2 and \mathbb{R}^3 are separate.



[†] Kerr et al., SIAM J. Sci. Comput. 30 (6) (2009) 3126

- **Definition of the problem:**

To capture the behavior of particles **with spatial extent** based on their rigid body features, dynamics, diffusion and hydrodynamic interactions.

- Coarse-graining the structure of molecules: Representing them as rigid bodies. Depending upon their symmetries,
 - Spherical and (a)symmetric tops in \mathbb{R}^3 .
 - Elliptic or circular thin discs in \mathbb{R}^2 .

- **How can we accomplish our goals (without a major change in the code base)?**

- Choice of integrators for 2d and 3d cases: Efficient, low relative error regarding the deterministic part.
- Choice of the overlap algorithm for the collision detection: Efficient, reusable in 2d and 3d cases with less effort (not require to implement separate analytical geometric methods).
- Computationally inexpensive way to do the scattering calculations during collision without effecting the concurrent stochastic events.

The geometric integrator for Langevin systems

- The Langevin type equations in the form of Ito of rigid bodies under the influence of conservative forces, hydrodynamic interactions, and thermal noise[†]:
 - Linear motion

$$\begin{aligned}dR^i &= \frac{p^i}{m^i} dt \\dP^i &= f^i(R, Q)dt - \sum_{j=1}^n \xi_{tt}^{ij}(R, Q) \frac{p^i}{m^j} dt \\&\quad - \frac{1}{2} \sum_{j=1}^n \xi_{tr}^{ij}(R, Q) A^T(Q^j) D^j \hat{S}^T(Q^j) \Pi^j dt \\&\quad - \frac{1}{2} \sum_{j=1}^n b_{tt}^{ij}(R, Q) dW^j + \sum_{j=1}^n b_{tr}^{ij}(R, Q) dW^j\end{aligned}$$

where n is the number of rigid bodies, $R^i(0) = r^i$ and $P^i(0) = p^i$ are the center of mass coordinates and corresponding conjugate momenta for rigid body i , respectively.

$$b(r, q)b^T(r, q) = 2k_B T \xi(r, q)$$

[†]Davidchack et al. J. Chem. Phys. 2017.

The geometric integrator for Langevin systems

- Rotational motion

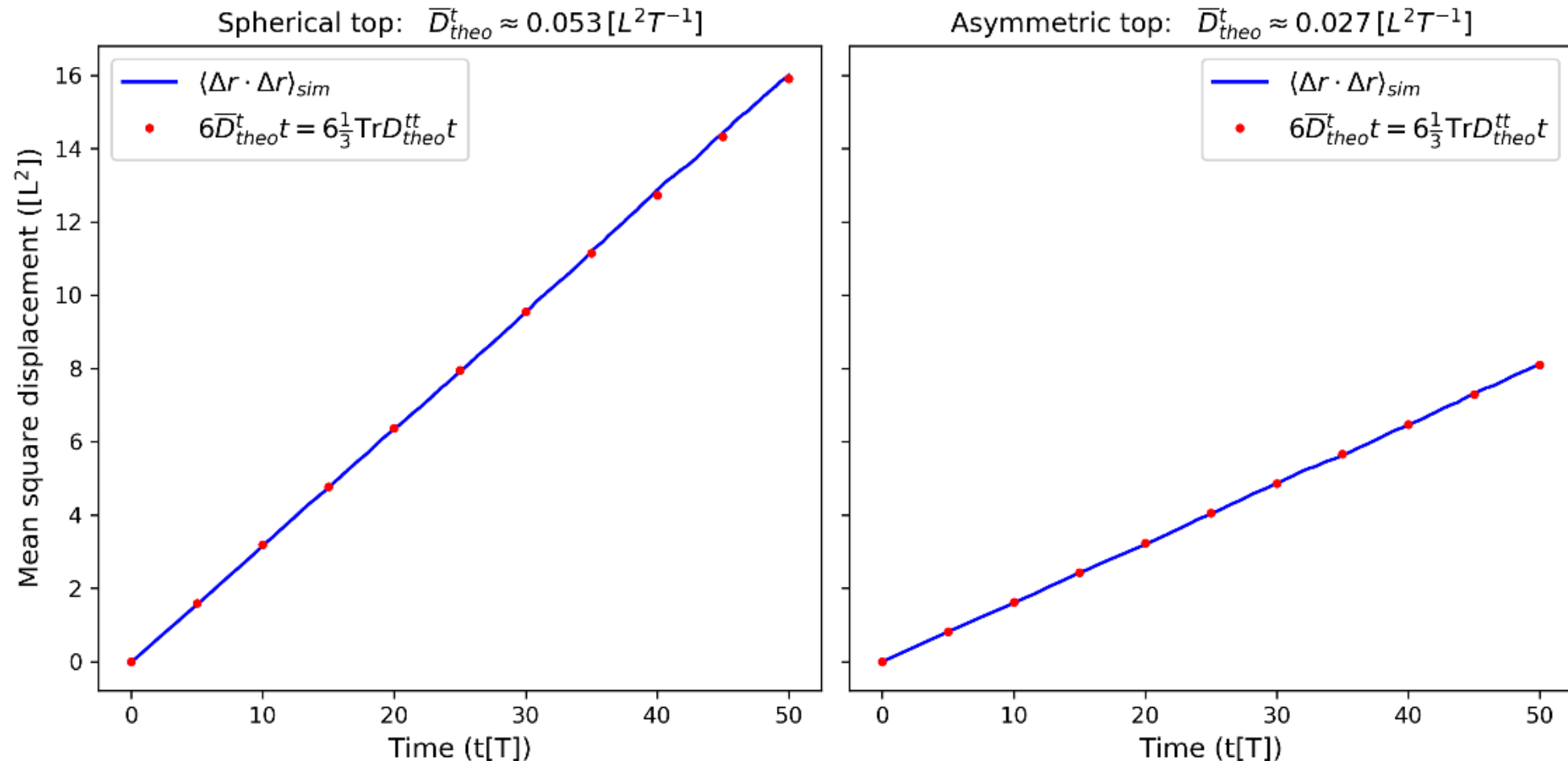
$$\begin{aligned}
 dQ^i &= \frac{1}{4} \hat{S}(Q^i) D^i \hat{S}^T(Q^i) \Pi^i dt \\
 d\Pi^i &= F^i(R, Q) dt + \frac{1}{4} \hat{S}(\Pi^i) D^i \hat{S}^T(Q^i) \Pi^i dt \\
 &\quad - \sum_{j=1}^n \check{S}(Q^i) \xi_{rr}^{ij}(R, Q) A^T(Q^j) D^j \hat{S}(Q^j) \Pi^j dt \\
 &\quad - 2 \sum_{j=1}^n \check{S}(Q^i) \xi_{rt}^{ij}(R, Q) \frac{P^j}{m^j} dt \\
 &\quad + 2 \sum_{j=1}^n \check{S}(Q^i) b_{rr}^{ij}(R, Q) dW^j + 2 \sum_{j=1}^n \check{S}(Q) b_{tr}^{ij}(R, Q) dW^j
 \end{aligned}$$

where $Q^i(0) = q^i$ and $\Pi^i(0) = \pi^i$ are the quaternions and the corresponding conjugate momenta, respectively.

$$q^{iT} \pi^i = 0, \quad |q^i| = 1$$

Simulation results without any external forces

Parallely generated ensemble of 10,000 independent trajectories for both a spherical and an asymmetric top. The simulation times and number of steps are chosen as 50 [T] and 10^4 respectively. Since the integrator has a natural over-damped limit, both bodies translationally diffuse with the overall average translational self-diffusion constants along their principal axes.



Collision Detection

- We should be careful when dealing with collision of bodies in stochastic processes.
- Stochastic kicks instantaneously change both the center of mass positions and the orientations of the bodies about to collide.
- We need two phases regarding the numerical integration:
 - Stochastic phase
 - Deterministic phase: Stochastic variables within the time interval Δt are assumed to be frozen, hence deterministic collision could be applicable.
- Continuously turning-on and -off the stochastic part of the integrator depending on the numerical integration scheme with the correct placement of the overlap algorithm!
- Possible issues:
 - Retroactive collision detection could be expensive (time bisection).
 - In 2d case, estimates of instantaneous translational and rotational pre-velocities while using overdamped equations could result in erroneous net impulse transfer, thereby effecting post-velocities and related coordinates.
 - Careless implementation of collision can result in increasing errors during long simulations, thus leading to a poor thermodynamic limit. In MCell, long simulation time is a sought-after feature.

Toy model in \mathbb{R}^2 - a rough collision detection

- 3 equivalent **point-particles** represented as circular discs to visualize a roughly estimated collision.
- Numerical integration scheme: First order point-particle overdamped integrator.
- Directions of their instantaneous velocities are estimated in terms of the changes in their coordinate vectors.
- Their speeds are assigned as the same and equal to the most probable speed w.r.t. 2d Maxwell-Boltzmann velocity distribution.
- Whenever an overlap is approximately detected within a predetermined threshold in the next frame, we calculate the impulse transferred during the collision to obtain the post-velocities. Then new coordinates are determined by those.
- Detected collisions are displayed with a “!!! BAM !!!” effect in the following video. Number of time steps are also shown during the animation.



Efficient overlap algorithm for ellipsoids

- An efficient overlap algorithm for rigid bodies should be implemented in two phases:
 - Broad-phase: No need to run an overlap check if the c.o.m. distances between a pair of bodies are larger than a spatial cutoff, i.e. they do not sense each other.
 - Narrow-phase: They are close enough to have a possible collision, therefore an overlap algorithm should be triggered for all such pairs.
- Perram-Wertheim overlap algorithm[†] is the ideal one for hard ellipsoids/ellipses. It is efficient and precise. It can be applied to determine possible overlaps in 2d and 3d cases. It can potentially provide
 - Contact points on the ellipsoids.
 - An upper bound to estimate the shortest directional distance between two ellipsoids' possible contact points measured along the direction of interbody vector connecting their c.o.m.'s.

[†]Perram et al. J. Comp. Phys. 58 (1985) 409.

Perram-Wertheim† Overlap Algorithm

Let's define two ellipsoids by their quadratic forms:

$$A(x) = (x - r)A(x - r), \quad B(x) = (x - s)B(x - s),$$
$$A = \sum_i a_i^{-2} u_i \otimes u_i, \quad B = \sum_i b_i^{-2} v_i \otimes v_i,$$

where u_i, v_i are the semi-axis vectors and a_i, b_i are their lengths. Let's define the following affine combination

$$S(\lambda, x) = \lambda A(x) + (1 - \lambda)B(x),$$

where $\lambda \in [0,1]$. The solution of the following optimization problem

$$F(A, B) = \max_{\lambda} \min_x S(\lambda, x)$$

allows us to determine three possible mutual configurations of the ellipsoids:

$$F(A, B) \begin{cases} < 1 \\ = 1 \\ > 1 \end{cases} \text{ for A and B } \begin{cases} \text{overlapping,} \\ \text{externally tangential,} \\ \text{nonoverlapping.} \end{cases}$$

†Perram et al. J. Comp. Phys. 58 (1985) 409.

Application to diffusing rigid bodies represented by elliptic thin discs in \mathbb{R}^2

- 3 equivalent 2d rigid bodies represented by elliptic discs.
- Diffusions are governed by the overdamped Langevin equations in the lab frame.
- Random directions on the discs stand for possible binding sites.
- The outliers “:”, “--” and “being filled” represent broad-phase (long-distance), non-overlapping in narrow-phase (intermediate distance) and overlapping cases (when we need to implement the retroactive collision detection), respectively.
- Time steps are shown in the animation.



Thank you for your patience!

Collaborators:

Tom Bartol, Ivet Bahar, Terrence Sejnowski